

# HODGE STRUCTURES AND WEIERSTRASS $\sigma$ -FUNCTION

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ABSTRACT. In this paper we introduce new definition of Hodge structures and show that  $\mathbb{R}$ -Hodge structures are determined by  $\mathbb{R}$ -linear operators that are annihilated by the Weierstrass  $\sigma$ -function

## 1. INTRODUCTION

Classically a Hodge structure of a given weight can be defined in the four equivalent ways as follows (see e.g. [2], [5]):

**Definition 1.1.** A Hodge structure of a weight  $n$  on a real vector space  $V$  consists of a finite-dimensional  $\mathbb{R}$ -vector space  $V$  together with any of the following equivalent data:

- (i) A decomposition  $V_{\mathbb{C}} = \bigoplus_{p+q=n} V^{p,q}$ , called the *Hodge decomposition*, such that  $\overline{V^{p,q}} = V_{q,p}$ .
- (ii) A decreasing filtration  $F_H^r V_{\mathbb{C}}$  of  $V_{\mathbb{C}}$ , called the *Hodge filtration*, such that  $F_H^r V_{\mathbb{C}} \oplus \overline{V_{\mathbb{C}}^{n-r+1}} = V_{\mathbb{C}}$ .
- (iii) A homomorphism  $h_n : \mathbb{S} \rightarrow \mathrm{GL}(V_{\mathbb{R}})$  of real algebraic groups, and also specifying that the weight of the Hodge structure is  $n$ .
- (iv) A homomorphism  $h_n : \mathbb{S} \rightarrow \mathrm{GL}(V_{\mathbb{R}})$  of real algebraic groups such that via the decomposition  $\mathbb{G}_m/\mathbb{R} \rightarrow \mathbb{S} \rightarrow \mathrm{GL}(V_{\mathbb{R}})$  an element  $t \in \mathbb{G}_m/\mathbb{R}$  acts as  $t^{-n} \cdot \mathrm{Id}$ .

Throughout the paper we work with Hodge structures of various weights, hence by a Hodge structure we understand here a finite direct sum

$$(1) \quad \rho := \bigoplus_{j=1}^k h_{n_j}$$

of representations  $h_{n_j}$  described in (iii) or (iv) of the Definition 1.1.

In this paper we consider Hodge structures on real vector space  $V$  via representations of the Lie algebra of the real algebraic group  $\mathbb{S}$  (denoted also  $\mathbb{C}^\times$ ) on  $V$ . In section 2 we show that a Hodge structure can be treated as a pair of operators  $E, T$  on  $V$  satisfying certain conditions (see Theorem 2.1). In section 3 we show that a Hodge structure can be treated as a single operator  $S := E + T$  on  $V$  such that  $\sigma(S) = 0$  for a Weierstrass  $\sigma$ -function which corresponds to decomposition of  $V$  into eigenspaces of operators  $E$  and  $T$ . Weierstrass  $\sigma$ -function does not have multiple zeros hence this corresponds to the fact that complexification of  $S$  does not have generalized eigenvectors other than ordinary ones.

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## 2. HODGE STRUCTURES AND LIE ALGEBRAS.

The following theorem gives another definition of the Hodge structure.

**Theorem 2.1.** *Let  $V$  be a finite dimensional vector space over  $\mathbb{R}$ . There is a one to one correspondence between the family of Hodge structures on  $V$  and the family of pairs of endomorphisms  $E, T \in \text{End}_{\mathbb{R}}(V)$  satisfying the following conditions:*

$$(2) \quad [E, T] = 0, \quad \sin(\pi E) = 0, \quad \sinh(\pi T) = 0,$$

$$(3) \quad \sin\left(\frac{\pi}{2}(E^2 + T^2)\right) = 0$$

*Proof.* Consider a Hodge structure on  $V$ . By (1) (cf. Definition 1.1 (iii)) this gives a representation:

$$\rho : \mathbb{S} \rightarrow \text{GL}(V)$$

of real algebraic groups. The representation  $\rho$  decomposes into irreducible representations  $\rho_{p,q}$  with multiplicities  $m_{p,q}$

$$\begin{aligned} \rho &= \bigoplus_{q \leq p} m_{p,q} \rho_{p,q}, \\ \rho_{p,q}(re^{i\phi}) &:= r^{p+q} \begin{bmatrix} \cos(p-q)\phi & -\sin(p-q)\phi \\ \sin(p-q)\phi & \cos(p-q)\phi \end{bmatrix} \quad \text{for } p \neq q, p, q \in \mathbb{Z} \\ \rho_{p,p}(re^{i\phi}) &:= r^{2p} \begin{bmatrix} 1 & \\ & 1 \end{bmatrix}. \end{aligned}$$

Certainly, the complexification of the representation  $\rho_{p,q}$  for  $q < p$  decomposes into two one-dimensional  $\mathbb{C}$ -vector spaces:

$$(4) \quad \rho_{p,q} \otimes_{\mathbb{R}} \mathbb{C} = \rho_{p,q}^{\mathbb{C}} \oplus \rho_{q,p}^{\mathbb{C}},$$

where

$$(5) \quad \rho_{m,n}^{\mathbb{C}}(z) = z^m \bar{z}^n.$$

Consider the real Lie algebra representation (the derivative of  $\rho$ ):

$$\mathcal{L}(\rho) : \mathbb{C} \rightarrow \text{End}(V).$$

For  $q \leq p$  the representation  $\mathcal{L}(\rho_{p,q})$  is also two-dimensional

$$\mathcal{L}(\rho_{p,q})(1) = (p+q)I \quad \text{and} \quad \mathcal{L}(\rho_{p,q})(i) = (p-q)J,$$

where

$$I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad J = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$$

For  $p = q$

$$\mathcal{L}(\rho_{p,p})(1) = 2p \quad \text{and} \quad \mathcal{L}(\rho_{p,p})(i) = 0.$$

If we put  $E := \mathcal{L}(\rho)(1)$  and  $T := \mathcal{L}(\rho)(i)$  then we get equations (2) and (3). The condition (3) is fulfilled because  $p - q$  and  $p + q$  have the same parity.

Now let us assume that conditions (2) and (3) hold. Observe that  $\sinh(z)$  and  $\sin(z)$  have single zeros in the complex plane. Moreover (2) and (3) imply that the complexifications  $E \otimes 1$  and  $T \otimes 1 \in \text{End}(V \otimes_{\mathbb{R}} \mathbb{C})$  have common eigenbasis. From this it follows that the endomorphisms  $E, T \in \text{End}_{\mathbb{R}}(V)$  have common Jordan decomposition into eigenspaces of dimension 1 or 2. We define a representation

$$\rho : \mathbb{C}^{\times} \rightarrow \text{GL}(V),$$

$$\rho(e^{x+iy}) = \exp(xE + yT) \text{ for } x, y \in \mathbb{R}.$$

$\rho$  is an algebraic representation, because the equality (3) holds. The representation  $\rho$  gives the Hodge structure on  $V$ .  $\square$

### 3. HODGE STRUCTURES VIA SINGLE OPERATOR

Let  $\sigma(z)$  be the Weierstrass' sigma function for the lattice generated by  $\omega_1 = 1 - i$  and  $\omega_2 = 1 + i$ :

$$\sigma(z) := z \prod_{(k_1, k_2) \neq (0,0)} \left( 1 - \frac{z}{k_1\omega_1 + k_2\omega_2} \right) \exp \left[ \frac{z}{k_1\omega_1 + k_2\omega_2} + \frac{1}{2} \left( \frac{z}{k_1\omega_1 + k_2\omega_2} \right)^2 \right]$$

**Theorem 3.1.** *For operators  $E, T \in \text{End}_{\mathbb{R}}(V)$  considered above let  $S := E + T$ . We get the following equality*

$$(6) \quad \sigma(S) = 0.$$

*Conversely every  $S \in \text{End}_{\mathbb{R}}(V)$  satisfying condition (6) gives a unique pair  $(E, T)$  of operators in  $\text{End}_{\mathbb{R}}(V)$  such that  $S = E + T$  and the conditions (2) and (3) hold.*

*Proof.* It is clear that  $S = E + T$  satisfies the equation (6). Conversely, assume that an operator  $S \in \text{End}_{\mathbb{R}}(V)$  satisfies (6). Since the  $\sigma$  function has zeros of order 1, we observe that the complexification of  $S$  is diagonalizable. We get the operators  $E$  and  $T$  considering equation

$$(7) \quad S(v) = \lambda v$$

in the complexification of  $V$ . The eigenvalues have integer real and imaginary parts with the same parity:

$$(8) \quad \lambda = a + ib, \quad a, b \in \mathbf{Z}, \quad a - b \in 2\mathbf{Z}.$$

Moreover we define the operators  $E, T$  in such a way that their complexifications acting on the eigenvector  $v$  of  $S$  have form:  $E(v) = av$  i  $T(v) = ibv$  where  $S(v) = (a + ib)v$ . Operators  $E$  and  $T$  satisfy equations (2) and (3). The operators  $E$  and  $T$  are uniquely determined. Indeed, if  $S = E' + T'$  such that  $E'$  and  $T'$  satisfy (2) and (3) then it is clear that  $[E', S] = 0$  and  $[T', S] = 0$ .  $\square$

**Remark 3.2.** For certain Hodge structures the set of eigenvalues of the complexification of  $S$  has further obstructions beyond (8). In this case  $S$  satisfies the equation  $g(S) = 0$ , where  $g(z)$  is an analytic function that divides  $\sigma(z)$  in such a way that  $\frac{\sigma(z)}{g(z)}$  is also an analytic function on the whole complex plane.

**Remark 3.3.** In our work in progress we define certain deformations of Hodge structures that arise in a natural way in mathematical physics (see [1], [3], [4]).

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